



# Measurement of the $W^+W^-$ and $ZZ$ production cross sections in pp collisions at $\sqrt{s} = 8$ TeV <sup>☆</sup>

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## ABSTRACT

The  $W^+W^-$  and  $ZZ$  production cross sections are measured in proton–proton collisions at  $\sqrt{s} = 8$  TeV with the CMS experiment at the LHC in data samples corresponding to an integrated luminosity of up to  $5.3 \text{ fb}^{-1}$ . The measurements are performed in the leptonic decay modes  $W^+W^- \rightarrow \ell'\nu\ell''\nu$  and  $ZZ \rightarrow 2\ell 2\ell'$ , where  $\ell = e, \mu$  and  $\ell'(\ell'') = e, \mu, \tau$ . The measured cross sections  $\sigma(\text{pp} \rightarrow W^+W^-) = 69.9 \pm 2.8(\text{stat.}) \pm 5.6(\text{syst.}) \pm 3.1(\text{lum.}) \text{ pb}$  and  $\sigma(\text{pp} \rightarrow ZZ) = 8.4 \pm 1.0(\text{stat.}) \pm 0.7(\text{syst.}) \pm 0.4(\text{lum.}) \text{ pb}$ , for both Z bosons produced in the mass region  $60 < m_Z < 120 \text{ GeV}$ , are consistent with standard model predictions. These are the first measurements of the diboson production cross sections at  $\sqrt{s} = 8$  TeV.

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## 1. Introduction

The study of  $W^+W^-$  and  $ZZ$  production in proton–proton collisions provides an important test of the standard model (SM). Any deviations of the measured cross sections from SM predictions would indicate new physics. Measurements of electroweak  $W^+W^-$  and  $ZZ$  production are essential for an accurate estimate of irreducible backgrounds for Higgs boson studies.

Previous measurements of  $W^+W^-$  and  $ZZ$  production were performed at the Large Hadron Collider (LHC) at a center-of-mass energy  $\sqrt{s} = 7$  TeV. With a data set corresponding to an integrated luminosity of  $36 \text{ pb}^{-1}$ , the Compact Muon Solenoid (CMS) Collaboration measured the  $W^+W^-$  cross section  $\sigma(\text{pp} \rightarrow W^+W^-) = 41.1 \pm 15.3(\text{stat.}) \pm 5.8(\text{syst.}) \pm 4.5(\text{lum.}) \text{ pb}$  [1], in good agreement with the SM prediction of  $47 \pm 2 \text{ pb}$  from Ref. [2]. The ATLAS Collaboration measured  $\sigma(\text{pp} \rightarrow W^+W^-) = 51.9 \pm 2.0(\text{stat.}) \pm 3.9(\text{syst.}) \pm 2.0(\text{lum.}) \text{ pb}$  [3] using  $4.6 \text{ fb}^{-1}$  of data. The  $ZZ$  cross section measurement from CMS used  $5 \text{ fb}^{-1}$  of data; the measured value,  $\sigma(\text{pp} \rightarrow ZZ) = 6.24^{+0.86}_{-0.80}(\text{stat.})^{+0.41}_{-0.32}(\text{syst.}) \pm 0.14(\text{lum.}) \text{ pb}$ , is consistent with the SM prediction of  $6.3 \pm 0.4 \text{ pb}$  for both Z bosons in the mass range  $60 < m_Z < 120 \text{ GeV}$  [4]. ATLAS measured  $\sigma(\text{pp} \rightarrow ZZ) = 6.7 \pm 0.7(\text{stat.})^{+0.4}_{-0.3}(\text{syst.}) \pm 0.3(\text{lum.}) \text{ pb}$  [5] with a data sample corresponding to an integrated luminosity of  $4.6 \text{ fb}^{-1}$ . Measurements of the  $W^+W^-$  and  $ZZ$  cross sections performed at the Tevatron are summarized in Refs. [6–10]. All measurements are found to agree well with the corresponding SM predictions.

In this Letter, the first measurements of the  $W^+W^-$  and  $ZZ$  production cross sections at  $\sqrt{s} = 8$  TeV are presented. The analysis is based on data collected in 2012 with the CMS experiment at the LHC, corresponding to an integrated luminosity of  $3.5 \text{ fb}^{-1}$  for the  $W^+W^-$  measurement and  $5.3 \text{ fb}^{-1}$  for the  $ZZ$  measurement. The measurements are performed in the  $W^+W^- \rightarrow \ell'\nu\ell''\nu$  and  $ZZ \rightarrow 2\ell 2\ell'$  decay channels, where  $\ell$  is e or  $\mu$ , and  $\ell'(\ell'')$  is e,  $\mu$ , or  $\tau$ . If a  $\tau$  lepton is present in the  $W^+W^-$  final state, only leptonic decays of the  $\tau$  lepton are considered. If a  $\tau$  lepton is present in the  $ZZ$  final state, one Z is required to decay either into  $e^+e^-$  or  $\mu^+\mu^-$ , and the second Z into  $\tau^+\tau^-$  in four possible final states:  $\tau_h\tau_h$ ,  $\tau_e\tau_h$ ,  $\tau_\mu\tau_h$ , and  $\tau_e\tau_\mu$ , where  $\tau_h$  indicates a  $\tau$  lepton decaying hadronically, and  $\tau_e$  and  $\tau_\mu$  indicate taus decaying into an electron and a muon, respectively.

The SM background sources to the  $W^+W^-$  event sample include  $W\gamma^{(*)}$ , top-quark ( $t\bar{t}$  and  $tW$ ),  $Z/\gamma^* \rightarrow \ell^+\ell^-$ , and diboson ( $WZ$  and  $ZZ$ ) production, as well as  $W$  + jets and QCD multijet events, where at least one of the jets is misidentified as a lepton. The SM background sources to the  $ZZ$  event sample include contributions from  $Zb\bar{b}$  and  $t\bar{t}$  processes, where the final states contain two isolated leptons and two b jets with secondary leptons, and from  $Z$  + jets and  $ZW$  + jets processes where the jets are misidentified as leptons.

## 2. The CMS detector and simulation

While the CMS detector is described in detail elsewhere [11], the key components for this analysis are summarized here. The CMS experiment uses a right-handed coordinate system, with the origin at the nominal interaction point, the  $x$  axis pointing to the

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center of the LHC ring, the  $y$  axis pointing up (perpendicular to the plane of the LHC ring), and the  $z$  axis along the anticlockwise-beam direction. The polar angle  $\theta$  is measured from the positive  $z$  axis and the azimuthal angle  $\phi$  is measured in the  $x$ - $y$  plane. The magnitude of the transverse momentum ( $p_T$ ) is calculated as  $p_T = \sqrt{p_x^2 + p_y^2}$ . A superconducting solenoid occupies the central region of the CMS detector, providing an axial magnetic field of 3.8 T parallel to the beam direction. The silicon pixel and strip tracker, the crystal electromagnetic calorimeter, and the brass/scintillator hadron calorimeter are located within the solenoid. A quartz-fibre Cherenkov calorimeter extends the coverage to  $|\eta| < 5.0$ , where pseudorapidity is defined as  $\eta = -\ln[\tan(\theta/2)]$ . Muons are measured in gas ionization detectors embedded in the steel flux return yoke outside the solenoid. The first level of the CMS trigger system, composed of custom hardware processors, is designed to select the most interesting events in less than 3  $\mu$ s using information from the calorimeters and muon detectors. The high-level-trigger processor farm decreases the event rate from 100 kHz delivered by the first level trigger to a few hundred hertz, before data storage.

Several Monte Carlo (MC) event generators are used to simulate the signals and backgrounds. The  $W^+W^-$  production via  $q\bar{q}$  annihilation is generated with the MADGRAPH [12] event generator, and PYTHIA [13] is used for parton showering, hadronization, and the underlying event simulation. The  $gg \rightarrow W^+W^-$  process, which is expected to contribute 3% of the total  $W^+W^-$  production rate [14], is generated with GG2WW [15]. The ZZ production via  $q\bar{q}$  annihilation is generated at next-to-leading order (NLO) with POWHEG 2.0 [16–18]. The  $gg \rightarrow ZZ$  process is simulated with GG2ZZ [19]. Other diboson processes ( $WZ$ ,  $W\gamma^{(*)}$ ) and the Z+jets samples are generated with MADGRAPH. The  $t\bar{t}$  and  $tW$  events are generated at NLO with POWHEG. For leading-order (LO) generators, the default set of parton distribution functions (PDF) used to produce these samples is CTEQ6L [20], while CT10 [21] is used for NLO generators. The  $\tau$  lepton decays are generated with TAUOLA [22]. For all processes, the detector response is simulated using a detailed description of the CMS detector, based on the GEANT4 package [23], and event reconstruction is performed with the same algorithms as used for data. The simulated samples include additional interactions per bunch crossing (pileup). The simulated events are weighted so that the pileup distribution matches the data, with an average pileup of about 20 interactions per bunch crossing.

### 3. Event reconstruction

Both the  $W^+W^-$  and ZZ event selections begin with the reconstruction and identification of lepton candidates. Electrons are reconstructed by combining information from the electromagnetic calorimeter and tracker [24,25]. Their identification relies on a multivariate technique that combines observables sensitive to the amount of bremsstrahlung along the electron trajectory, the geometrical and momentum matching between the electron trajectory in the tracker and the energy deposit in the calorimeter, as well as the shower shape [24]. Muons are reconstructed [26] with information from both the tracker and the muon spectrometer, and are required to pass selection criteria similar to those described in Ref. [1]. A particle-flow (PF) technique [27] is used to reconstruct  $\tau_h$  candidates with the “hadron plus strip” (HPS) algorithm [28], which is designed to optimize the performance of  $\tau_h$  identification and reconstruction by considering specific  $\tau_h$  decay modes. In the PF approach, information from all subdetectors is combined to reconstruct and identify particles produced in the collision. The particles are classified into mutually exclusive categories: charged hadrons, photons, neutral hadrons, muons, and electrons. These particles are used to reconstruct the  $\tau_h$  candidates; the neutrinos

produced in all  $\tau$  decays escape detection and are ignored in the  $\tau_h$  reconstruction. The lepton candidates are required to be consistent with the primary vertex of the event, which is chosen as the vertex with highest  $\sum p_T^2$  of its associated tracks. This criterion provides the correct assignment for the primary vertex in more than 99% of both signal and background events for the pileup distribution observed in the data.

Charged leptons from W and Z boson decays are usually isolated from other activity in the event. For each electron or muon candidate, a cone is constructed around the track direction at the event vertex. The scalar sum of the transverse momenta of all reconstructed particles consistent with the chosen primary vertex and contained within the cone is calculated, excluding the contribution from the lepton candidate itself. To improve the discrimination against nonisolated muons from the  $W$ +jets background in the  $W^+W^-$  selection, this procedure is repeated with several cones of different widths. This isolation information is then combined by means of a multivariate technique. For both electrons and muons a correction is applied to account for the energy contribution in the isolation cone due to pileup. A median transverse energy due to pileup is determined event by event and is subtracted from the transverse energy ( $E_T$ ) in the isolation cone [29]. A similar technique is used to form  $\tau$  lepton isolation quantities.

Jets are reconstructed from the PF particles using the anti- $k_T$  clustering algorithm [30] with distance parameter of 0.5, as implemented in the FASTJET package [31,32]. The jet energy is corrected for pileup in a manner similar to the correction of the energy inside a lepton isolation cone. Jet energy corrections are also applied as a function of the jet  $p_T$  and  $\eta$  [33]. A multivariate selection is applied to separate jets coming from the primary interaction from those reconstructed using energy deposits associated with pileup. This discrimination is based on the differences in the jet shapes and in the relative multiplicity of charged and neutral components. Tracks associated with a jet are required to be consistent with the primary vertex.

To suppress the top-quark background in events without high- $p_T$  jets, top-quark-tagging techniques are defined with two methods. The first method vetoes events containing muons originating from b quarks [34] appearing in top-quark decays. The second method uses b-jet tagging applied to jets with  $15 < p_T < 30$  GeV based on tracks with large impact parameter within jets.

The missing transverse energy  $\vec{E}_T^{\text{miss}}$  is defined as the negative vector sum of the transverse momenta of all reconstructed particles in the event. A projected  $E_T^{\text{miss}}$  is defined as (i) the magnitude of the  $\vec{E}_T^{\text{miss}}$  component transverse to the closest lepton, if  $\Delta\phi(\ell, \vec{E}_T^{\text{miss}}) < \pi/2$ , or (ii) the magnitude of the  $\vec{E}_T^{\text{miss}}$  otherwise. This observable more efficiently rejects  $Z/\gamma^* \rightarrow \tau^+\tau^-$  background events in which the  $\vec{E}_T^{\text{miss}}$  is preferentially aligned with the leptons, and  $Z/\gamma^* \rightarrow \ell^+\ell^-$  events with mismeasured  $\vec{E}_T^{\text{miss}}$ . Since the projected  $E_T^{\text{miss}}$  resolution is degraded as pileup increases, the minimum of two different observables is used: the first includes all particle candidates in the event, while the second uses only the charged particle candidates associated with the primary vertex.

## 4. Event selection and background estimates

### 4.1. $W^+W^-$ production

Events are selected with two oppositely charged electron or muon candidates, both with  $p_T > 20$  GeV and with  $|\eta| < 2.5$  for the electrons and  $|\eta| < 2.4$  for the muons. The  $\tau$  leptons contribute to the measurement only if they decay to electrons or muons that pass the selection requirements. At the trigger level, events are required to have a pair of electrons or muons where one of the

leptons has  $p_T > 17$  GeV and the other  $p_T > 8$  GeV, or a single electron (muon) with  $p_T > 27$  (24) GeV. The trigger efficiency is approximately 98% for both  $q\bar{q} \rightarrow W^+W^-$  and  $gg \rightarrow W^+W^-$  processes.

To reduce the background from top-quark decays, events with one or more jets surviving the jet selection criteria and with  $p_T > 30$  GeV and  $|\eta| < 4.7$  are rejected. The residual top-quark background is further suppressed by 50% after applying the top-quark-tagging techniques.

In order to reduce the Drell–Yan background, the projected  $E_T^{\text{miss}}$  is required to be above 45 GeV in the  $e^+e^-$  and  $\mu^+\mu^-$  final states. For the  $e^\pm\mu^\mp$  final state, which has a lower contamination from  $Z/\gamma^* \rightarrow \ell^+\ell^-$  decays, the threshold is reduced to 20 GeV. These requirements remove more than 99% of the Drell–Yan background.

To further reduce the Drell–Yan background in the  $e^+e^-$  and  $\mu^+\mu^-$  final states, the angle in the transverse plane between the dilepton system total momentum and the most energetic jet with  $p_T > 15$  GeV is required to be smaller than 165 degrees. Events with dilepton masses within  $\pm 15$  GeV of the Z mass or below 12 GeV are also rejected. Finally, the transverse momentum of the dilepton system,  $p_T^{\ell\ell}$ , is required to be above 45 GeV to reduce contributions from Drell–Yan background and events containing jets misidentified as leptons.

To reduce the background from other diboson processes, such as WZ or ZZ production, any event is rejected if it has a third lepton with  $p_T > 10$  GeV passing the identification and isolation requirements. The  $W\gamma^{(*)}$  production in which the photon converts is suppressed by rejecting electrons consistent with a photon conversion.

The W + jets and QCD multijet backgrounds are estimated from a control region in which one lepton passes the nominal requirements, while the other passes looser criteria on impact parameter and isolation, but fails the nominal requirements. The contribution to the control region from processes with two genuine leptons is subtracted by using simulation. The number of events in the signal region is obtained by scaling the number of events in the control region with the efficiency for loosely identified lepton candidates to pass the tight selection. These efficiencies are measured in data using multijet events and are parametrized according to the  $p_T$  and  $|\eta|$  of the lepton candidate.

The normalization of the top-quark background that survives the top-quark-tagging requirements is estimated from data by counting the number of top-tagged events and applying the corresponding top-tagging efficiency. The top-tagging efficiency is measured with a control sample dominated by  $t\bar{t}$  and  $tW$  events, which is selected by requiring a b-tagged jet.

We estimate the Drell–Yan contribution to the  $e^+e^-$  and  $\mu^+\mu^-$  final states outside of the Z mass window by normalizing the event yield from simulation to the observed number of events inside the Z mass window. The methods used to estimate both the top-quark and Drell–Yan backgrounds are described in more detail in Ref. [35].

Finally, a control sample with three reconstructed leptons is used to measure the data-to-MC scaling factor for the  $W\gamma^{(*)}$  process. We use only  $W\gamma^{(*)} \rightarrow \ell\nu\mu^+\mu^-$  events for this measurement because the  $\ell\nu e^+e^-$  final state is difficult to separate from other backgrounds. This measurement is used to normalize the simulated  $W\gamma^{(*)}$  background contribution from asymmetric gamma decays in which one lepton escapes detection [36].

Other backgrounds, such as WZ and ZZ diboson production, are estimated from simulation.

The W + jets and  $W\gamma^{(*)}$  background estimate is checked using data events that pass all the selection requirements with the exception that the two leptons must have the same charge. After

subtraction of the expected WZ background, this sample is dominated by W + jets and  $W\gamma^{(*)}$  events. The  $Z/\gamma^* \rightarrow \tau^+\tau^-$  contamination is checked using  $Z/\gamma^* \rightarrow e^+e^-$  and  $Z/\gamma^* \rightarrow \mu^+\mu^-$  events selected in data, where the leptons are replaced with simulated  $\tau$  lepton decays.

#### 4.2. ZZ production

Selected events are required to have at least one electron or muon with  $p_T > 20$  GeV and another one with  $p_T > 10$  GeV, and  $|\eta| < 2.5$  (2.4) for electrons (muons). All other electrons (muons) are required to have  $p_T > 7$  (5) GeV. The  $\tau_h$  candidates are required to have  $p_T > 20$  GeV and  $|\eta| < 2.3$ . All leptons must originate from the same vertex and be isolated. At the trigger level, events are required to have a pair of electrons or muons, one lepton with  $p_T > 17$  GeV and the other with  $p_T > 8$  GeV.

The selected events are required to contain two Z candidates. One candidate, denoted by  $Z_1$ , should decay into electrons or muons,  $Z \rightarrow \ell^+\ell^-$ , and must have reconstructed invariant mass  $60 < m_{\ell\ell} < 120$  GeV. If more than one candidate is found, the one with mass closest to the Z mass is considered as  $Z_1$ .

The selection requirements for the second Z candidate, denoted by  $Z_2$ , depend on the final state. In the  $4\mu$ ,  $4e$ , and  $2e2\mu$  final states the isolation requirements are the same as for the leptons from  $Z_1$ . For the  $e^+e^-\tau_e^\pm\tau_\mu^\mp$  and  $\mu^+\mu^-\tau_e^\pm\tau_\mu^\mp$  final states the electron and muon  $p_T$  values are required to exceed 10 GeV. In final states with  $Z_2 \rightarrow \tau_e\tau_h, \tau_\mu\tau_h$  the isolation requirements for all the electrons and muons are tighter. A study of inclusive  $Z \rightarrow \tau^+\tau^-$  production [37] shows that modifying the electron and muon isolation requirements is a more effective way to reduce background in such final states than imposing tighter isolation criteria on  $\tau_h$ .

The invariant mass of the reconstructed  $Z_2$  is required to satisfy  $60 < m_{\ell\ell} < 120$  GeV when  $Z_2$  decays into  $e^+e^-$  or  $\mu^+\mu^-$ . In the  $2\ell 2\tau$  final states, the visible invariant mass of the reconstructed  $Z_2 \rightarrow \tau^+\tau^-$  is shifted to smaller values because of undetected neutrinos in  $\tau$  decays. Therefore, in the final states involving  $\tau$  leptons, the visible mass is required to satisfy  $30 < m_{\tau\tau} < 90$  GeV, and the leptons from the same Z are required to be separated by  $\Delta R > 0.4$  for the  $Z_1$ , and by  $\Delta R > 0.5$  for the  $Z_2$ .

Estimated acceptances of the selection requirements defined with respect to the full phase space are 58%, 56%, 54%, and 25% for the  $4e$ ,  $2e2\mu$ ,  $4\mu$ , and  $2\ell 2\tau$  final states, respectively.

The major contributions to the background come from Z production in association with jets, WZ production in association with jets, and  $t\bar{t}$ . In all these cases, a jet or nonisolated lepton is misidentified as an isolated electron, muon, or  $\tau_h$ . The relative contribution of each source of background depends on the final state.

For the background estimation, two different approaches are used. Both start by relaxing the isolation and identification criteria for two additional reconstructed lepton objects indicated as  $\ell_{\text{reco}}\ell_{\text{reco}}$  in the  $Z_1 + \ell_{\text{reco}}\ell_{\text{reco}}$  event sample. The additional pair of leptons is required to have like-sign charge (to avoid signal contamination) and same flavor ( $e^\pm e^\pm, \mu^\pm \mu^\pm$ ). The first method estimates the number of Z + X background events in the signal region by taking into account the lepton misidentification probability for each of the two additional leptons. The second method uses a control region with two opposite-sign leptons that fail the isolation and identification criteria. The background in the signal region is estimated by weighting the events in the control region with the lepton misidentification probability. In addition, a control region with three passing leptons and one failing lepton is used to account for contributions from backgrounds with three prompt leptons and one misidentified lepton. Comparable background rates

in the signal region are found within the uncertainties from both methods.

## 5. Systematic uncertainties

The uncertainty in the signal acceptance for the two measurements due to variations in the parton distribution functions and the value of  $\alpha_s$  is estimated by following the PDF4LHC prescription [38]. Using CT10 [21], MSTW08 [39], and NNPDF [40] sets, the uncertainties are 2.3% (4.0%) for the  $q\bar{q} \rightarrow W^+W^-$  (ZZ) processes.

The effects of higher-order corrections are found by varying the QCD renormalization and factorization scales simultaneously up and down by a factor of two using the mCFM program [14]. The variations in the acceptance are found to be 1.5% for the  $q\bar{q} \rightarrow W^+W^-$  process and to be negligible for the  $q\bar{q} \rightarrow ZZ$  and  $gg \rightarrow ZZ$  processes.

The  $W^+W^-$  jet veto efficiencies in data are estimated from simulation, and multiplied by a data-to-simulation scale factor derived from  $Z/\gamma^* \rightarrow \ell^+\ell^-$  events in the Z peak:  $\epsilon_{W^+W^-}^{\text{data}} = \epsilon_{W^+W^-}^{\text{MC}} \times \epsilon_Z^{\text{data}}/\epsilon_Z^{\text{MC}}$ . The uncertainty is thus factorized into the uncertainty in the Z efficiency in data and the uncertainty in the ratio of the  $W^+W^-$  efficiency to the Z efficiency in simulation ( $\epsilon_{W^+W^-}^{\text{MC}}/\epsilon_Z^{\text{MC}}$ ). The former, which is statistically dominated, is 0.3%. Theoretical uncertainties due to higher-order corrections contribute most to the  $W^+W^-/Z$  efficiency ratio uncertainty, which is estimated to be 4.6% for  $W^+W^-$  production. The data-to-simulation correction factor is found to be close to one and is not applied.

Simulated events are scaled according to the lepton efficiency correction factors measured using data control samples, which are typically close to one. The uncertainties in the measured identification and isolation efficiencies are found to be 1–2% for muons and electrons, and 6–7% for  $\tau_h$ . The uncertainty in the trigger efficiency is 1.5%. The uncertainty in the lepton energy scale is about 3% for  $\tau_h$ , and 1–2.5% and 1.5% for electrons and muons, respectively.

The uncertainties in the Z + jets, WZ + jets, and  $t\bar{t}$  backgrounds to the  $ZZ \rightarrow 2\ell 2\ell'$  selection are 30–50% depending on the decay channel. These uncertainties comprise the statistical and systematic uncertainties in the misidentification rates measured in data control samples.

The systematic uncertainties in the W + jets, Z + jets, and top backgrounds to the  $W^+W^- \rightarrow \ell'\nu\ell''\nu$  selection are 36%, 24%, and 15%, respectively. The theoretical uncertainties in the WZ and ZZ cross sections are calculated following the same prescription as for the signal acceptance. Including the experimental uncertainties gives a systematic uncertainty in WZ and ZZ backgrounds of approximately 10%.

The uncertainty assigned to the pileup reweighting procedure amounts to 2.3%. The uncertainty in the integrated luminosity is 4.4% [41].

## 6. Results

### 6.1. $W^+W^-$ cross section measurement

The observed and expected signal plus background yield is summarized in Table 1. The expected  $W^+W^-$  contribution is calculated assuming the SM cross section.

The total background yield is  $275 \pm 35$  events and the expected signal and background yield is  $959 \pm 60$  events, with 1111 events observed.

The measured  $W^+W^-$  yield is calculated by subtracting the estimated contributions of the various background processes. The product of the signal efficiency and acceptance averaged over all

**Table 1**

Expected and observed event yields for the  $W^+W^-$  selection. The uncertainties correspond to the statistical and systematic uncertainties added in quadrature.

Channel	$\ell'\nu\ell''\nu$
$W^+W^-$	$684 \pm 50$
$t\bar{t}$ and $tW$	$132 \pm 23$
W + jets	$60 \pm 22$
WZ and ZZ	$27 \pm 3$
Z/ $\gamma^*$ + jets	$43 \pm 12$
$W\gamma^{(*)}$	$14 \pm 5$
Total background	$275 \pm 35$
Signal + background	$959 \pm 60$
Data	1111

lepton flavors including  $\tau$  leptons is  $(3.2 \pm 0.2)\%$ . Using the  $W \rightarrow \ell\nu$  branching ratio of  $0.1080 \pm 0.0009$  from Ref. [42], the  $W^+W^-$  production cross section in pp collision data at  $\sqrt{s} = 8$  TeV is measured to be

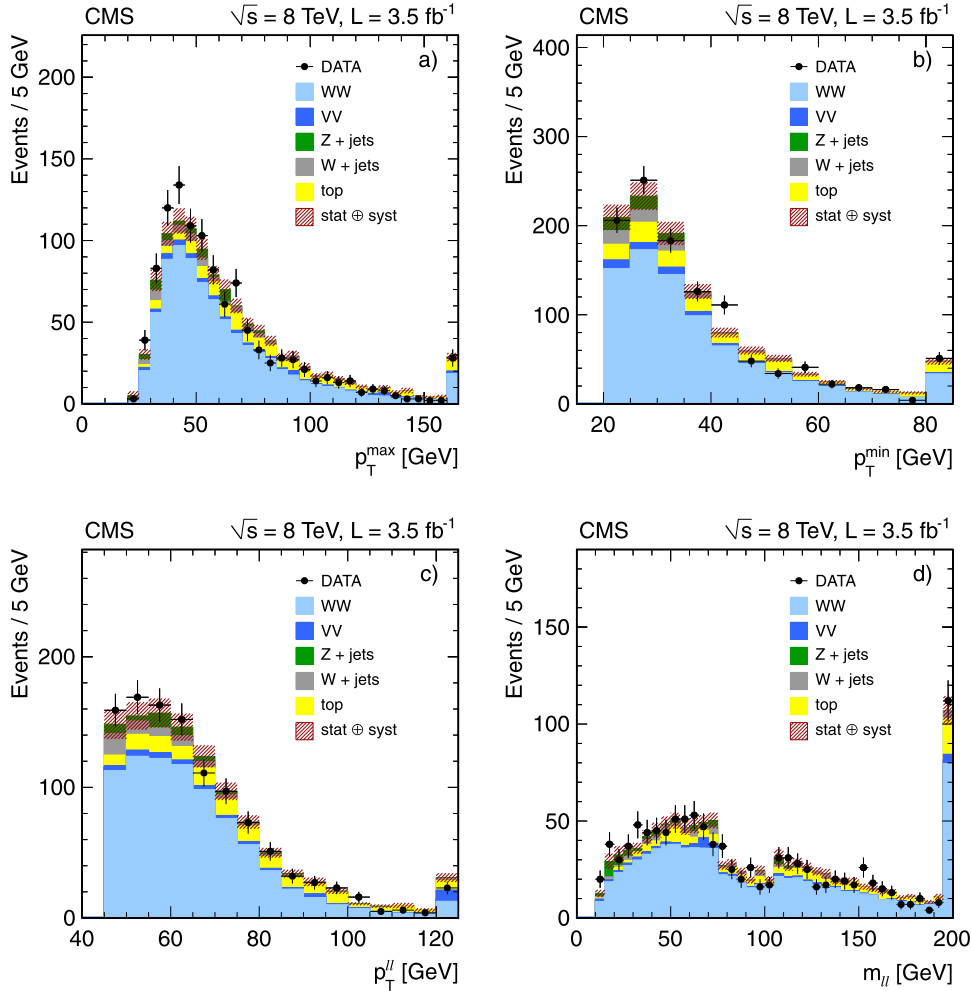
$$\begin{aligned} \sigma(\text{pp} \rightarrow W^+W^-) \\ = 69.9 \pm 2.8(\text{stat.}) \pm 5.6(\text{syst.}) \pm 3.1(\text{lum.}) \text{ pb.} \end{aligned}$$

The statistical uncertainty reflects the total number of observed events. The systematic uncertainty includes both the statistical and systematic uncertainties in the background prediction, as well as the uncertainty in the signal efficiency. This measurement is slightly higher than the SM expectation of  $57.3^{+2.3}_{-1.6}$  pb, calculated in Ref. [2] by using mCFM at NLO with the MSTW08 PDF and setting the factorization and renormalization scales to the W mass. Additional processes may increase the production yield in the  $W^+W^-$  final state by as much as 5% for the event selection used in this analysis. Higgs boson production would give an additional contribution of about 4% of the cross section given above, based on next-to-next-to-leading-order cross section calculations for the  $H \rightarrow W^+W^-$  process [43] under the assumption that the newly discovered resonance [44,45] is a SM-like Higgs boson with a mass of 125 GeV. Contributions from diffractive production [46], double parton scattering, and QED exclusive production [47] are also considered.

The distributions of the leading lepton transverse momentum  $p_T^{\text{max}}$ , the trailing lepton transverse momentum  $p_T^{\text{min}}$ , the dilepton transverse momentum  $p_T^{\ell\ell}$ , and the dilepton invariant mass  $m_{\ell\ell}$  are shown in Fig. 1, where the  $W^+W^-$  contribution is normalized to the measured cross section.

### 6.2. ZZ cross section measurement

Table 2 presents the observed and expected yields and the number of the estimated background events in the signal region. There are 71 candidates observed in the  $4e, 4\mu$ , and  $2e2\mu$  channels, to be compared to an expectation of  $65.6 \pm 4.4$  events. Among the expected events 1.4 are from background processes. In the  $2\ell 2\tau$  channels 13 candidates are observed. The expected  $12.1 \pm 1.6$  events for  $2\ell 2\tau$  channels contain 5.6 events from background processes. The reconstructed four-lepton invariant mass distributions are shown in Figs. 2(a) and 2(b) for the sum of the  $4e, 4\mu$ , and  $2e2\mu$  channels, and the sum of all the  $2\ell 2\tau$  channels. Data are compared to the SM expectations. The shapes of the signal and the background are taken from the MC simulation, with each component normalized to the corresponding estimated value from Table 2. The reconstructed masses in  $2\ell 2\tau$  states are



**Fig. 1.** Distributions for  $W^+W^-$  candidate events of (a) the leading lepton transverse momentum  $p_T^{\max}$ , (b) the trailing lepton transverse momentum  $p_T^{\min}$ , (c) the dilepton transverse momentum  $p_T^{\ell\ell}$ , and (d) the dilepton invariant mass  $m_{\ell\ell}$ . Points represent the data, and shaded histograms represent the  $W^+W^-$  signal and the background processes. The last bin includes the overflow. The  $W^+W^-$  signal is scaled to the measured cross section, and the background processes are normalized to the corresponding estimated values in Table 1.

**Table 2**

Expected and observed event yields for the ZZ selection. The uncertainties correspond to the statistical and systematic uncertainties added in quadrature.

Channel	4e	4 $\mu$	2e2 $\mu$	2 $\ell$ 2 $\tau$
ZZ	11.6 $\pm$ 1.4	20.3 $\pm$ 2.2	32.4 $\pm$ 3.5	6.5 $\pm$ 0.8
Background	0.4 $\pm$ 0.2	0.4 $\pm$ 0.3	0.5 $\pm$ 0.4	5.6 $\pm$ 1.4
Signal + background	12.0 $\pm$ 1.4	20.7 $\pm$ 2.2	32.9 $\pm$ 3.5	12.1 $\pm$ 1.6
Data	14	19	38	13

shifted downwards with respect to the generated values by about 30% because of the undetected neutrinos in  $\tau$  decays. Figs. 2(c) and 2(d) demonstrate the relationship between reconstructed  $Z_1$  and  $Z_2$  masses.

To measure the ZZ cross section the numbers of observed events are unfolded in a combined likelihood fit. Each decay mode is treated as a separate channel giving eleven measurements to combine: 4e, 4 $\mu$ , 2e2 $\mu$ , and eight 2 $\ell$ 2 $\tau$  channels. The  $\tau$ -lepton decay modes are treated separately; the methodology used for event reconstruction and selection ensures that the decay modes are mutually exclusive. The joint likelihood is a combination of the likelihoods for the individual channels, which include the signal and background hypotheses. The statistical and systematic uncertainties are introduced in the form of nuisance parameters via

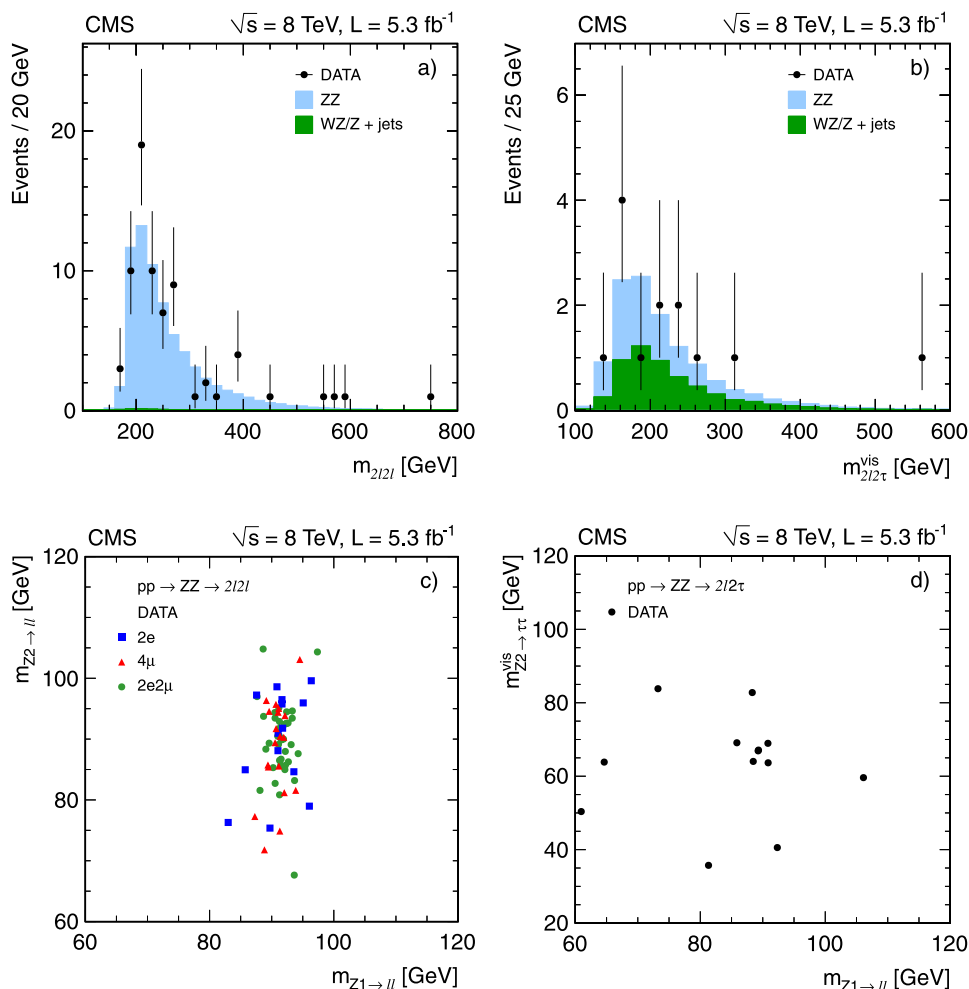
log-normal distributions around the estimated central values. The resulting cross section is

$$\begin{aligned} \sigma(\text{pp} \rightarrow \text{ZZ}) \\ = 8.4 \pm 1.0(\text{stat.}) \pm 0.7(\text{syst.}) \pm 0.4(\text{lum.}) \text{ pb.} \end{aligned}$$

This is to be compared to the theoretical value of  $7.7 \pm 0.4$  pb calculated with mCFM at NLO for  $q\bar{q} \rightarrow \text{ZZ}$  and LO for  $gg \rightarrow \text{ZZ}$  with MSTW08 PDF, and factorization and renormalization scales set to the Z mass, for both lepton pairs in the mass range  $60 < m_Z < 120$  GeV.

## 7. Summary

The  $W^+W^-$  and ZZ production cross sections have been measured in proton–proton collisions at  $\sqrt{s} = 8$  TeV in the  $W^+W^- \rightarrow \ell'\nu\ell''\nu$  and  $\text{ZZ} \rightarrow 2\ell 2\ell'$  decay modes with  $\ell = e, \mu$  and  $\ell'(\ell'') = e, \mu, \tau$ . The data samples correspond to an integrated luminosity of  $3.5 \text{ fb}^{-1}$  for the  $W^+W^-$  and  $5.3 \text{ fb}^{-1}$  for the ZZ measurements. The measured production cross sections  $\sigma(\text{pp} \rightarrow W^+W^-) = 69.9 \pm 2.8(\text{stat.}) \pm 5.6(\text{syst.}) \pm 3.1(\text{lum.}) \text{ pb}$  and  $\sigma(\text{pp} \rightarrow \text{ZZ}) = 8.4 \pm 1.0(\text{stat.}) \pm 0.7(\text{syst.}) \pm 0.4(\text{lum.}) \text{ pb}$ , for both Z bosons produced in the mass region  $60 < m_Z < 120$  GeV, are consistent with the



**Fig. 2.** Distributions for ZZ candidate events of (a) the four-lepton reconstructed mass for the sum of the  $4e$ ,  $4\mu$ , and  $2e2\mu$  channels and (b) the sum of the  $2l2\tau$  channels. Points represent the data, and shaded histograms represent the expected ZZ signal and the reducible background. The shapes of the signal and background are taken from the MC simulation, with each component normalized to the corresponding estimated value from Table 2. The distributions (c) and (d) demonstrate the relationship between the reconstructed  $Z_1$  and  $Z_2$  masses. Different symbols are used to present different decay channels.

standard model predictions. This is the first measurement of the diboson production cross sections at  $\sqrt{s} = 8$  TeV.

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